### 1 Analogue closed-loop optogenetic modulation of hippocampal pyramidal cells

### 2 dissociates gamma frequency and amplitude

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#### 12 Abstract

13 Gamma-band oscillations are implicated in modulation of attention, integration of sensory

14 information and flexible communication among anatomically connected brain areas. How

15 networks become entrained is incompletely understood. Specifically, it is unclear how the

16 spectral and temporal characteristics of network oscillations can be altered on rapid

17 timescales needed for efficient communication. We use closed-loop optogenetic modulation

18 of principal cell excitability in mouse hippocampal slices to interrogate the dynamical

19 properties of hippocampal oscillations. Gamma frequency and amplitude can be modulated

20 bi-directionally, and dissociated, by phase-advancing or delaying optogenetic feedback to

21 pyramidal cells. Closed-loop modulation alters the synchrony rather than average frequency

22 of action potentials, in principle avoiding disruption of population rate-coding of information.

23 Modulation of phasic excitatory currents in principal neurons is sufficient to manipulate

24 oscillations, suggesting that feed-forward excitation of pyramidal cells has an important role

25 in determining oscillatory dynamics and the ability of networks to couple with one another.

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### 27 Impact statement

28 Neurons can synchronize, supporting flexible communication among brain areas; closed-

29 loop optogenetics allows the frequency and power of population oscillations to be

- 30 dissociated, providing a tool to interrogate how networks couple.
- 31

#### 32 Introduction

33 Gamma-band (approximately 30 to 120 Hz) oscillations have been implicated in the 34 modulation of attention and perception, in action initiation, spatial navigation and memory 35 encoding, and have also been proposed to underlie flexible information routing among 36 anatomically connected regions (Akam & Kullmann, 2010, 2014; Börgers & Kopell, 2003; 37 Fries, 2005; Kirst, Timme, & Battaglia, 2016; Lisman, 2010; Rodriguez et al., 1999; Salinas & 38 Sejnowski, 2001; Schnitzler & Gross, 2005). Central to several of these proposed roles is the ability of gamma oscillations in different areas to enter into, and exit, states of synchrony with 39 40 one another (Akam, Oren, Mantoan, Ferenczi, & Kullmann, 2012; Fries, 2015; Varela, 41 Lachaux, Rodriguez, & Martinerie, 2001). Evidence for behavioural-state dependent coupling 42 and uncoupling comes from variable oscillatory coherence among distinct components of the 43 visual cortex, correlating with selective stimulus attention (Bosman et al., 2012; Grothe, 44 Neitzel, Mandon, & Kreiter, 2012). An earlier study in the rodent hippocampal formation 45 showed that the CA1 subfield can flip between a state of coherence with the medial 46 entorhinal cortex at ~110 Hz and a state of coherence with the CA3 subfield at ~40 Hz, 47 correlating with information flow through the temporo-ammonic and Schaffer collateral 48 pathways respectively (Colgin et al., 2009). Although several experimental confounds cloud 49 the interpretation of coherence measured from local field potential (LFP) recordings (Buzsáki 50 & Schomburg, 2015), these studies provide some of the most compelling evidence that 51 gamma-band oscillatory entrainment underlies flexible functional connectivity.

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53 Although the cellular mechanisms underlying gamma oscillations have been extensively 54 studied (Bartos, Vida, & Jonas, 2007; Buzsáki & Wang, 2012), there remain uncertainties 55 over the fundamental determinants of their dynamics and the relative contributions of 56 excitatory and inhibitory signalling. Gamma-band oscillations can be induced in vitro in the 57 presence of blockers of ionotropic glutamate receptors (Whittington, Traub, & Jefferys, 58 1995), or in vivo by optogenetic stimulation of parvalbumin-positive interneurons (Cardin et 59 al., 2009; Sohal, Zhang, Yizhar, & Deisseroth, 2009), underlining the importance of fast 60 perisomatic inhibition (Bartos et al., 2002; Fisahn et al., 2004; Mann, Radcliffe, & Paulsen, 61 2005). Robust population oscillations can also be simulated in exclusively inhibitory networks 62 (Wang & Buzsáki, 1996). These experimental and computational observations emphasize 63 the importance of inhibitory kinetics. Nevertheless, gamma-band oscillations can be 64 entrained by sinusoidal optogenetic stimulation of pyramidal neurons in an in vitro 65 hippocampal slice preparation (Akam et al., 2012). This observation implies that phasic depolarization of principal cells can determine the gamma rhythm and argues against a 66 67 model where the only role of pyramidal cells is to tonically depolarize a network of 68 reciprocally coupled interneurons (Bartos et al., 2007; Tiesinga & Sejnowski, 2009).

70 Further insight into the dynamical mechanisms of synchronization between oscillating 71 networks comes from examining the phase response curve (PRC) of the network oscillation, 72 defined as the phase advance or delay produced by a transient stimulation, as a function of 73 the instantaneous phase at which the stimulus is delivered. The finding that gamma in an in 74 vitro hippocampal slice preparation shows a biphasic PRC (Akam et al., 2012) is consistent 75 with the hypothesis that this oscillation can be entrained by appropriately modulated afferent 76 activity. The shape of the PRC is furthermore accurately reproduced with a simple neural 77 mass model (Wilson & Cowan, 1972), where extracellular electrical or optogenetic stimuli 78 are represented as transient perturbations of the instantaneous level of excitation or 79 inhibition (Akam et al., 2012). Recent theoretical work has derived population PRCs for 80 oscillations in spiking network models, providing an insight into how mechanisms of 81 oscillation generation determine entrainment properties (Akao, Ogawa, Jimbo, Ermentrout, & 82 Kotani, 2018; Kotani, Yamaguchi, Yoshida, Jimbo, & Ermentrout, 2014). Nevertheless, there 83 remains a large gap between the PRC and understanding the determinants of the oscillatory 84 frequency and interactions between gamma-generating circuits.

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86 The present study investigates the dynamical properties of gamma oscillations by using 87 closed-loop optogenetics to create an artificial feedback loop between the oscillatory network 88 activity (as assessed by the LFP) and excitatory input to the principal cell population. 89 Specifically, we delivered analogue-modulated excitation whose strength was a function of 90 the instantaneous phase and amplitude of the oscillation. This approach is guite distinct from 91 previous closed-loop applications of optogenetics (Grosenick, Marshel, & Deisseroth, 2015), 92 which have adopted one of four main strategies. First, several studies have used the 93 detection of a change in the state of a network, such as the onset of an electrographic 94 seizure (Krook-Magnuson, Armstrong, Oijala, & Soltesz, 2013; Paz et al., 2013) or sharp-95 wave ripple (Stark et al., 2014), to trigger light delivery and return the network to its ground 96 state. Second, light pulses have been timed according to the phase of a theta oscillation 97 (Siegle & Wilson, 2014), while examining the consequences for behaviour. In the latter 98 example the theta oscillation itself was not altered. Third, optogenetics has been used to 99 regulate the overall activity of a population of neurons at a desired level (Newman et al., 100 2015). Fourth, optogenetic depolarization of interneurons, triggered by spikes in an individual 101 principal cell, has been used to simulate a feedback inhibitory loop to interrogate their role in 102 gamma (Sohal et al., 2009; Veit, Hakim, Jadi, Sejnowski, & Adesnik, 2017). The goal of the 103 present investigation is qualitatively different: to understand how the spectral characteristics 104 of gamma are affected by rhythmic excitation arriving at different phases. Computational 105 simulations have suggested that closed loop optogenetics could be used to adjust the phase

106	of gamma (Witt et al., 2013), but whether it can alter its frequency or amplitude remains
107	unclear.

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#### 111 Results

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#### 113 Closed-loop feedback modulation of affects gamma oscillations in CA1

114 We expressed the red-shifted optogenetic actuator C1V1 (Yizhar et al., 2011) in the mouse

115 hippocampus CA1 under the *Camk2a* promoter to bias expression to excitatory neurons.

116 The local field potential (LFP) was recorded in the CA1 pyramidal cell layer in acute

117 hippocampal slices. A slowly increasing ramp of light (peak wavelength 590 nm) was

delivered via a light-emitting diode (LED) coupled to the epifluorescence port of an upright

microscope, eliciting a gamma oscillation (**Fig. 1a, b**), as previously reported in rodents

120 (Adesnik, 2018; Adesnik & Scanziani, 2010; Akam et al., 2012; Butler, Mendonça, Robinson,

4 Paulsen, 2016; Pastoll, Solanka, van Rossum, & Nolan, 2013), cats (Ni et al., 2016) and

122 monkeys (Lu et al., 2015).

123 In order to investigate the role of phasic excitation in setting the dynamical properties of 124 gamma we used the LFP itself to manipulate the optogenetic drive in real time. The LED 125 driver command was multiplied by a simple function of the instantaneous value of the LFP 126 and its time-derivative:  $(1 + k_1 LFP + k_2 dLFP/dt))$ , where  $k_1$  and  $k_2$  are positive or negative 127 constants. These operations were implemented with a field-programmable gate array 128 (FPGA) and applied for a defined duration (typically 1 or 2 seconds) during the ramp. This 129 yielded a change in the spectral properties of the oscillation, which lasted for the duration of 130 the closed-loop feedback (Fig. 1c). Because both the LFP and its time-derivative fluctuated 131 about 0, the "gamma clamp" had little effect on the average illumination intensity relative to 132 an unmodulated ramp. Changes in the oscillation frequency or power could therefore not be 133 attributed to a net increase or decrease in the average optogenetic drive to pyramidal 134 neurons.

135 We adjusted the clamp function by altering the values of  $k_1$  and  $k_2$  and asked whether the 136 frequency and/or power of the gamma oscillation can be modulated bidirectionally. Changes 137 in spectral properties were related to the phase difference between the LFP and the LED 138 drive during the clamp, as estimated from the cross-spectrum at maximal magnitude. In-139 phase modulation, achieved by setting  $k_1$  positive and  $k_2 = 0$ , led to an increase in oscillatory

140 power and frequency (Fig. 2a). Modulating the ramp in anti-phase relative to the LFP, by

- setting  $k_1$  negative, led to a decrease in both frequency and power (Fig. 2b). Advancing the
- 142 phase of the clamp by approximately 90°, achieved by setting  $k_1 = 0$  and  $k_2$  positive,
- 143 increased the frequency of the oscillation whilst decreasing is power (Fig. 2c). Finally, a
- 144 decrease in frequency and increase in power was achieved by delaying the trough of the
- 145 clamp modulation relative to the LFP, by setting  $k_2$  negative (**Fig. 2D**). Detailed inspection of
- 146 the ramp command waveform during the clamp shows that it was in some cases distorted
- relative to the LFP, as expected from its non-sinusoidal shape (Cole & Voytek, 2017) (e.g.
- 148 **Fig. 2c, d**), and so the LED-LFP phase differences were only approximate.
- 149 Attempts to estimate the instantaneous oscillation phase, for instance using a Hilbert
- 150 transform, and to use this to phase-advance or phase-delay a template of the LFP,
- 151 compressed or stretched in time, were unsuccessful: the phase jitter and cycle-to-cycle
- 152 variability in the amplitude and frequency of the gamma oscillation (see LFP traces in Fig. 2)
- 153 prevented accurate estimation of these parameters in the face of closed loop feedback.
- 154

#### 155 Oscillation clamp is broadly consistent with the phase response curve of gamma

- 156 Changes in frequency and power, expressed in relation to the approximate phase difference 157 between the LED command and the LFP, were qualitatively consistent across experiments 158 (**Fig. 3a–c**). Moreover, as the LED-LFP phase difference was rotated through a complete 159 cycle, the effect on the oscillation in the two-dimensional plane defined by the change in 160 oscillation frequency and power also rotated through 360°, such that with the appropriate 161 phase of closed-loop feedback the network oscillation could be pushed in any desired
- 162 direction in the oscillation frequency-power space (**Fig. 3c**).
- 163 To gain a mechanistic insight, we asked if the characteristic relationship between the 164 frequency change and the LED-LFP phase difference could be explained by the shape of 165 the phase-response curve (PRC) previously reported (Akam et al., 2012). In that study, a 166 brief 'kick' was applied on top of the LED ramp command, and the phase advance or delay 167 of subsequent oscillations was related to the phase of the LFP at which the transient 168 occurred. A phase delay was observed when the transient optogenetic stimulus was 169 delivered at the trough of the LFP, when pyramidal neurons are most likely to fire. The 170 maximal phase advance, in contrast, occurred when the stimulus was delivered 171 approximately one third of a cycle after the trough of the LFP. Assuming linear behaviour, 172 the effect of modulating the light intensity in closed loop can be obtained by averaging the 173 product of the phase shift and the LFP over the entire cycle of the oscillation. The circular 174 cross-correlogram between the typical LFP shape and the PRC should then predict the

effect of modulating the optogenetic drive by the shape of the LFP itself at arbitrary degrees of phase advance or delay (**Fig. 3d**). In-phase modulation is expected, on the basis of this calculation, to phase-advance the oscillation, and thus to result in an increase in oscillatory frequency over successive cycles. Anti-phase modulation, in contrast, is predicted to phasedelay the oscillation, and thus to decrease is frequency. The circular cross-correlation is, moreover, asymmetrical, broadly consistent with the shape of the relationship between the change in frequency and LED-LFP phase difference observed in the clamp experiments

182 (**Fig. 3a**).

183 Although the shape of the PRC is consistent with the changes in gamma frequency achieved 184 with closed loop modulation at different LED-LFP phase differences, on its own it says 185 nothing about changes in power. Power was maximally decreased with a phase advance of 186 the LED command over the LFP around 90°, whilst it was maximally increased with a phase 187 delay around 90° (Fig. 3b). The relative phases at which frequency and power were altered 188 are however consistent with the behaviour of a normal form description of a super-critical 189 Hopf bifurcation in the vicinity of its limit-cycle. In this scenario, the LFP would approximate 190 an observed variable, and the optogenetic drive would act in the direction of a hidden 191 variable at a  $+90^{\circ}$  angle to the LFP.

192 A deeper understanding of the characteristic changes in frequency and power of the LFP

193 with different LED-LFP phase differences require an insight into how neurons spike and

194 ultimately how membrane currents respond to the fluctuations in optogenetic drive. We

195 therefore performed single-cell recordings in parallel with the LFP recordings.

196

#### 197 Gamma clamp affects the timing, not rate, of pyramidal neuron firing

198 Although the average illumination intensity was not altered during the gamma clamp, for

199 certain LED-LFP phase relationships gamma power increased or decreased robustly.

200 Inhibitory currents in principal neurons, rather than spikes or excitatory currents, have

201 previously been shown to be the main determinant of the LFP (Oren, Hájos, & Paulsen,

202 2010), suggesting that the change in power during the clamp is not a direct effect of the

203 optogenetic drive but results instead from a change in pyramidal neuron synchrony or phase,

in a reciprocal relationship with the degree and temporal synchrony of interneuron

205 recruitment. To determine how the clamp affects pyramidal neuron firing, we repeated

206 experiments with an additional patch pipette to record from individual pyramidal neurons in

207 cell-attached mode. Individual action potentials were used to align the simultaneously

208 recorded LFP, and to estimate the phase at which they occurred. During an unmodulated

- ramp, pyramidal cells tended to spike sparsely, close to the trough of the oscillation,
- 210 consistent with previous studies of pharmacologically induced oscillations (Fisahn, Pike,
- Buhl, & Paulsen, 1998). During the clamp, an increase in oscillatory power was associated
- 212 with a corresponding increase in the degree of synchrony of pyramidal cell firing: the circular
- 213 dispersion of LFP phase at which pyramidal cells fired decreased relative to the unclamped
- situation (Fig. 4a). Conversely, a decrease in power was accompanied by a relative
- 215 desynchronization of pyramidal cell firing. This relationship was qualitatively consistent, as
- indicated by the change in vector length obtained from the circular average of spike phases
- 217 (Fig. 4c). The vector length increased when oscillation power increased (p = 0.02, n = 12,
- sign test), and decreased when oscillation power decreased (p = 0.025, n = 5). Strikingly,
- 219 however, there was no change in the overall firing rate of pyramidal cells when the oscillation
- 220 power was increased or decreased by the clamp. Changes in power were thus achieved by
- tightening the synchrony of firing, or by desynchronizing action potentials, rather than by
- altering the overall activity of pyramidal neurons.
- Increases in oscillatory frequency were accompanied by a phase advance of pyramidal cell firing relative to the LFP (**Fig. 4b, c,**  $p = 8 \times 10^{-7}$ , n = 7, Hotellier test (Zar, 2009)). A trend for a phase delay was observed in a small number of experiments where frequency-lowering clamp was tested (n = 3). This observation is consistent with the view that changes in the phase of pyramidal neuron action potentials are causally upstream of changes in gamma frequency, even though the current generators of the LFP itself are dominated by GABAergic signalling (Gulyás et al., 2010; Hájos et al., 2004; Oren et al., 2010).
- 230

# Excitatory current phase in principal cells determines changes in gamma spectralproperties

- In the examples illustrated in **Figs. 4a** and **b**, the optogenetic modulation was applied with a phase advance over the LFP of  $\sim 0^{\circ}$  and  $\sim 45^{\circ}$  respectively. Why does in-phase modulation
- result in an increase in power, and phase-advanced modulation result in an increase in
- frequency? To gain a mechanistic insight into how gamma clamp operates, we examined the
- 237 phase of excitation experienced by pyramidal neurons during different clamp regimes.
- 238 We repeated experiments as above, but with one pipette used to voltage–clamp a pyramidal
- 239 neuron at the estimated GABA<sub>A</sub> reversal potential (approximately –70 mV), and the other
- 240 pipette to record the LFP. We then measured the inward current at each phase of the
- gamma oscillation, as defined by the LFP, and repeated this over consecutive cycles to
- obtain an average time-course (**Fig. 5a**). The minimum (that is, least negative) inward

243 current during the average cycle was subtracted to yield an estimate of the phasic excitatory 244 current, which could then be represented as a vector representing its average phase and 245 amplitude (Fig. 5b). During unclamped gamma, the excitatory current was small, and its 246 average phase relative to the LFP varied among experiments, as expected from the very 247 sparse synaptic connectivity among pyramidal neurons in CA1 (Deuchars & Thomson, 248 1996). Gamma clamp imposed a large phasic inward current (Fig. 5b). Subtracting the 249 vector representing the baseline phasic inward current yielded a vector representing the net 250 excitatory current imposed by the gamma clamp ( $\Delta E$ ). This lagged behind the LED 251 modulation, reflecting in part the opsin activation and deactivation kinetics (Fig. 5c, and 252 arrows in Fig. 5b, right). For the example illustrated in Fig. 5 an 83° phase advance of the 253 LED over the LFP resulted in  $\Delta E$  with mean phase of 247°, where the LFP trough is defined 254 as 0°. This vielded an increase in frequency and decrease in power of the gamma

255 oscillation.

256 Comparing across different clamp regimes reveals how gamma frequency and power 257 change in relation to the phasic excitation experienced by principal cells (Fig. 6a, b). An 258 increase in gamma frequency was achieved when the average excitatory current phase 259 occurred during the down-stroke of the LFP (~180° to 360°), whilst a decrease in frequency 260 was achieved when excitation was applied during the upstroke (~0° to 180°). An increase in 261 power, on the other hand, was achieved with excitation around the trough ( $\sim 270^{\circ}$  to  $90^{\circ}$ ), 262 and a decrease in power occurred with excitation around the peak (~90° to 270°). Given 263 that, under baseline conditions, pyramidal neurons fire maximally near to the trough of the 264 LFP (0°), these data imply that the increase in frequency occurs because they are brought to 265 firing threshold earlier (see also **Fig. 4b, c**). An increase in power, on the other hand, occurs 266 because pyramidal neurons are synchronized by adding a depolarization when they are 267 most likely to fire (see also Fig. 4a, c).

268

#### **Gamma clamp affects inhibitory currents**

Finally, we asked how inhibitory currents in a subset of pyramidal neurons are altered by the
closed-loop optogenetic manipulation, by voltage-clamping them around the glutamate
reversal potential (0 mV). The mean phase and amplitude of outward inhibitory currents
were calculated in a similar way, by subtracting the minimal current from the circular average
of the outward GABA<sub>A</sub> receptor-mediated current (Fig. 7). In contrast to excitatory currents,
phasic inhibitory currents under baseline conditions were large, consistent with the major
role of feedback interneurons in gamma (Fig. 7a, b). Changes in inhibitory currents (ΔI) were

277 relatively smaller than for excitatory currents and were dominated by effects on the power of 278 the oscillation. Thus, for a +83° LED-LFP phase advance, which led to a decrease in gamma 279 power and increase in frequency (same cell as in Fig. 5), there was little change in the 280 average phase of the inhibitory current, although it was decreased (Fig. 7c). Aligning the 281 average currents by the LFP across different trials whilst the cell was held at -70 mV or 0 282 mV, and during unclamped and clamped periods, yielded an insight into the relationship 283 between excitatory and inhibitory conductances during the oscillatory cycle (Fig. 7d). For the 284 example shown in Fig. 5 and Fig. 7a-c, the effective cycle changed from one dominated by 285 phasic inhibition to one with similar relative amplitudes of phasic inhibition and excitation, 286 with excitation leading inhibition in both cases (Fig. 7d, left). In contrast, for a 133° LED-LFP 287 phase delay, which led to a decrease in frequency and increase in power, the inhibitory 288 current increased (Fig. 7d, right). Similar results were obtained in 8 cells.

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290

#### 291 Discussion

The present study shows that closed-loop optogenetic manipulation of principal cells allows predictable, bidirectional and dissociable changes in the power and frequency of gamma

294 oscillations. We observed a broad consistency between the frequency manipulation

achieved with closed loop optogenetic feedback and that predicted from the phase response

behaviour previously observed with intermittent optogenetic stimuli (Akam et al., 2012).

297 Optogenetically and pharmacologically induced gamma also exhibited similar dynamical

298 properties in that study, implying that the principles uncovered in the present work are not

specific to the way gamma oscillations were elicited.

300 Previous studies have stressed the importance of fast-spiking parvalbumin-positive (PV+) 301 interneurons in gamma (Cardin et al., 2009; Sohal et al., 2009) (but see (Veit et al., 2017)). 302 PV+ basket cells tend to fire with very little phase dispersion, close to one-to-one with each 303 cycle of the oscillation in vitro (Bartos et al., 2007; Gulyás et al., 2010). Our attempts to 304 achieve gamma clamp by targeting interneurons rather than pyramidal cells have thus far been unsuccessful because their out of phase recruitment powerfully suppresses the 305 306 oscillation (data not shown). The weaker phase-locking of pyramidal than PV+ cell firing to 307 gamma oscillations, together with their sparse firing on successive cycles of gamma 308 (Csicsvari, Jamieson, Wise, & Buzsáki, 2003; Gulyás et al., 2010; Tukker, Fuentealba, 309 Hartwich, Somogyi, & Klausberger, 2007), may however confer a broader dynamic range 310 over which they can influence the phase, frequency and amplitude of the oscillation. Taken 311 together with previous evidence that open-loop sinusoidal optogenetic stimulation of

312 principal cells can entrain a gamma oscillation (Akam et al., 2012), the present data

- 313 underline the importance of action potential timing in principal cells in the spectral and
- 314 temporal properties of hippocampal gamma, notwithstanding the evidence that the LFP itself
- is dominated by inhibitory currents in principal cells (Oren et al., 2010), and argue against a
- 316 model where the function of principal cells is only to depolarize a population of reciprocally
- 317 connected interneurons.

318 Closed-loop manipulations have been applied previously in the context of network

- 319 oscillations, using either electrical and optogenetic stimuli delivered at specific phases of
- 320 theta or gamma oscillations, in order to probe the mechanisms of long-term plasticity
- induction (Huerta & Lisman, 1995; Pavlides, Greenstein, Grudman, & Winson, 1988) or
- 322 sharp-wave ripple generation (Stark et al., 2014), or to test the theta phase-dependence of
- 323 memory encoding and retrieval (Siegle & Wilson, 2014). A similar strategy has been used to
- 324 interrupt experimental thalamocortical seizures (Berényi, Belluscio, Mao, & Buzsáki, 2012).
- 325 However, these studies have not aimed at modulating the amplitude or frequency of an on-
- 326 going oscillation.
- 327 We have focused on gamma because a local circuit is sufficient to generate the oscillation,
- 328 and we have previously shown that the phase response behaviour of hippocampal gamma is
- 329 well described by a simple dynamical model (Akam et al., 2012). The circuits underlying
- 330 theta and other oscillations either involve longer-range connections in the brain or are poorly
- defined. They are therefore less likely to be amenable to local optogenetic manipulation.
- 332 This does not exclude the possibility that, for instance, theta oscillations in the hippocampus
- 333 could be manipulated by closed-loop modulation of excitability in the basal forebrain.
- 334 The ability to alter the amplitude and frequency of gamma suggests a versatile tool to test 335 the roles of gamma in information routing and other high-level brain functions, both in health 336 and in disease states such as schizophrenia (Uhlhaas & Singer, 2010). Hitherto, most 337 experimental manipulations of oscillations have relied on periodic stimulation, which can 338 entrain network oscillations (Akam et al., 2012) or evoke oscillations in an otherwise 339 asynchronous network (Cardin et al., 2009; Sohal et al., 2009). Transcranial stimulation 340 designed to entrain oscillations in vivo can bias perception (Neuling, Rach, Wagner, Wolters, 341 & Herrmann, 2012; Romei, Gross, & Thut, 2010; Thut, Schyns, & Gross, 2011) and 342 bidirectionally affect performance in motor (Joundi, Jenkinson, Brittain, Aziz, & Brown, 2012) 343 and working memory (Polanía, Nitsche, Korman, Batsikadze, & Paulus, 2012) tasks. 344 However, external periodic stimulation is not well suited to desynchronize network activity or 345 to suppress oscillatory dynamics. Furthermore, if periodic stimulation is used, the desired 346 change in amplitude or frequency is achieved at the cost of imposing an externally

- determined phase on the oscillation. This will prevent the oscillation from entraining to
   endogenous periodic signals such as those arising from other oscillating networks or
   periodic sensory stimuli.
- 350 Closed-loop stimulation, in which signals recorded from a network are used in real time to 351 bias its state, in principle provides an alternative way of manipulating network oscillations, 352 and has been used to interfere with pathological rhythms in models of Parkinson's disease 353 (Rosin et al., 2011), to suppress Parkinsonian tremor (Brittain, Probert-Smith, Aziz, & Brown, 354 2013), and in a model of thalamocortical epilepsy (Butt et al., 2005). This approach relies on 355 an artificial feedback loop which either counteracts or amplifies the endogenous feedback 356 responsible for synchronizing the network (Rosenblum & Pikovsky, 2004). Importantly, 357 optogenetics has the advantage over electrical stimulation that the modulation can be 358 distributed across a population of neurons. We have, moreover, shown that closed-loop 359 manipulation of a gamma oscillation can be achieved without a net increase or decrease in 360 the average firing rate of neurons, implying that it would not necessarily perturb information 361 represented as an average firing rate code.
- 362 Extrapolating from *in vitro* gamma to the brain *in situ* presents several technical challenges,
- 363 including the need for optical fibers to illuminate the tissue and the potential for
- 364 photoelectrical artifacts. Moreover, oscillations are generally less prominent because the
- 365 current generators from multiple oscillating and non-oscillating populations overlap,
- 366 complicating the evaluation of phase and frequency. Nevertheless, the present study
- 367 identifies some general principles to guide attempts to achieve bidirectional and dissociable
- 368 modulation of oscillatory frequency and power *in vivo*. This should allow a definitive test of
- the causal role of gamma in functions such as attention modulation and information routing(Sohal, 2016).
- 371

#### 372 Materials & Methods

#### 373 Key resource table

Reagent type (species) or resource	Designation	Source or reference	Identifiers	Additional information
Strain, strain background ( <i>Mus</i>	Wild type, C57bl6 mice	UCL Biological		

Reagent type (species) or resource	Designation	Source or reference	Identifiers	Additional information
musculus)		Services		
Genetic reagent	Viral vector (C1V1): AAV5-CaMKIIa- C1V1(E122T/E162T)- TS-eYFP	UNC Vector Core	Serotype: 5 C1V1: AAV-CaMKIIa- C1V1(E122T/E162T)-TS-EYFP	
Other	Equipment (viral injections): Sterotaxic frame	Kopf Instruments	Model 900	
Chemical compound, drug	Salts, drugs	Sigma		
Other	Equipment (LED light source)	Cairn Instruments	OptoLED, 590 nm	
Other	Equipment (LED light source) Custom assembled LED	Thorlabs	High-power mounted LED: M590L2; Tube lens: SM1V10; Planoconvex lens: LA1951-A- N-BK7; Coupler: SM1T2; Adapter for microscope: SM1A14	
Other	Equipment (LED driver)	Thorlabs	DC2100	
Other	Equipment (Upright microscope)	Olympus	BX51WI, UMPLFLN 20X W IR	
Other	Equipment (Upright microscope)	Scientifica	SliceScope, UMPLFLN 20X W IR	
Other	Equipment (Amplifier)	Molecular Devices	Multiclamp 700B	
Other	Equipment (Data acquisition card)	National Instruments	PCI-6221	

Reagent type (species) or resource	Designation	Source or reference	Identifiers	Additional information
Other	Equipment (computing) Real-time controller	National Instruments	cRIO-9022	
Other	Equipment (computing). FPGA	Xilinx	Virtek-5	
Software, algorithm	LabVIEW	National Instruments	LabVIEW, LabVIEW Real-Time 2013, 2014, 2015, 2016, 2017	Custom virtual instruments
Software, algorithm	R	www.R- project.org	Ver. 3.3.0 for Mac	

- 376 All procedures followed the Animals (Scientific Procedures) Act, 1986, and were reviewed by 377 the UCL Institute of Neurology Animal Welfare and Ethical Review Body. P21 male C57 378 mice were anesthetized with isoflurane and placed in a stereotaxic frame (Kopf Instruments). 379 A suspension of AAV5-CaMKIIa-C1V1(E122T/E162T)-TS-eYFP (UNC Vector Core, titre 5 x 380 10<sup>12</sup> IU/ml) was injected at a rate of 100 nl/min into 4 sites in both hippocampi (injection 381 volume: 300-500 nl per site). The antero-posterior injection coordinate was taken as 2/3 of 382 the distance from bregma to lambda. The lateral coordinates were 3.0 mm from the midline, 383 and the ventral coordinates were 3.5, 3.0, 2.5 and 2.0 mm from the surface of the skull. 384 Hippocampal slices were prepared at least 4 weeks later. Animals were sacrificed by 385 pentobarbitone overdose and underwent transcardiac perfusion with an oxygenated solution 386 containing (in mM): 92 N-methyl-D-glucamine-Cl, 2.5 KCl, 1.25 NaH<sub>2</sub>PO<sub>4</sub>, 20 HEPES, 30 387 NaHCO<sub>3</sub>, 25 glucose, 10 MgCl<sub>2</sub>, 0.5 CaCl<sub>2</sub>, 2 thiourea, 5 Na-ascorbate and 3 Na-pyruvate, 388 with sucrose added to achieve an osmolality of 315 mOsm/L. Brain slices (400 µm thick) 389 were prepared at room temperature and then incubated at 37 °C for 12 minutes in the same 390 solution. They were subsequently stored at room temperature, in a solution containing (in 391 mM): 126 NaCl, 3 KCl, 1.25 NaH<sub>2</sub>PO<sub>4</sub>, 2 MgSO<sub>4</sub>, 2 CaCl<sub>2</sub>, 24 NaHCO<sub>3</sub>, 10 glucose, shielded 392 from light, before being transferred to the stage of an upright microscope (Olympus BX51WI 393 or Scientifica SliceScope), where they were perfused on both sides with the same solution at
- 394 32° C. Expression of C1V1 in CA1 was verified by epifluorescence, and CA3 was ablated to
- 395 focus on local gamma-generating mechanisms.
- Epifluorescence imaging and C1V1 stimulation were achieved with LEDs (OptoLED, Cairn
  Instruments, or assembled from Thorlabs components using an M590L2 590 nm LED and a
  DC2100 high-power LED driver). The light source was coupled to the epifluorescence
  illuminator of the microscope, with a silver mirror in the place of a dichroic cube. Wide-field
  illumination was delivered via a 20x, 0.5 NA water immersion objective. The current
  delivered to the LED was kept in the linear input-output range, and the irradiance was <5</li>
  mW/mm<sup>2</sup>. Light ramps typically lasting 8 s were delivered every 30 45 s.
- 403 LFPs were recorded in the CA1 pyramidal layer using patch pipettes filled with extracellular 404 solution and a Multiclamp 700B amplifier (Molecular Devices), and band-pass filtered 405 between 1 and 200 or 500 Hz. A linear LED ramp command was generated via a 406 multifunction data acquisition card (National Instruments PCI-6221) and, together with the 407 LFP, was digitized using a real-time controller (National Instruments cRIO-9022) with a Xilinx 408 Virtex-5 FPGA (cRIO-9133) operating at a loop rate of 10 kHz. The ramp was multiplied by 409  $(1 + k_1 \text{LFP} + k_2 d \text{LFP}/dt)$ , stepping through different values of *k* in a pseudo-random order for

- 410 successive trials. *d*LFP/*d*t was calculated as the difference between successive digitization
- 411 values in the FPGA, averaged over successive 2 ms intervals to minimize high-frequency
- 412 noise. The output of the FPGA/real-time controller was sent to the LED driver, and digitized
- 413 in parallel with the LFP at 10 kHz on the data acquisition PC.

414 To study the phase relationship of action potentials and the LFP oscillation, a cell-attached

415 recording was obtained using a second patch pipette held in voltage clamp mode, low-pass

416 filtered at 10 kHz and digitized in parallel with the LFP and LED command signal. The phasic

- 417 excitatory or inhibitory current was recorded in the same way, but using a whole-cell pipette
- 418 containing (in mM): K-gluconate (145), NaCl (8), KOH-HEPES (10), EGTA (0.2), Mg-ATP (2)
- 419 and Na 3 -GTP (0.3); pH 7.2; 290 mOsm. Phasic conductances (Fig. 7d) were estimated
- 420 from Ohm's law, assuming a driving force of 70 mV.

421 Off-line analysis was performed in LabVIEW (National Instruments) and R. Time-frequency

422 spectrograms were calculated using a Morlet wavelet transform and are displayed as heat

423 maps. Because the gamma oscillation was non-stationary, its frequency was estimated by

424 calculating the short-term Fourier transform and then averaging the mean instantaneous

425 frequency for successive overlapping intervals. The power of the oscillation was estimated in

426 the same way, by averaging the power at the mean instantaneous frequency.

427 Spikes were identified using threshold crossing. The instantaneous oscillation phase was

428 estimated by passing a 200-ms segment of the LFP centered on the spike through a

- 429 Hanning window, and then calculating its phase and frequency using the Extract Single Tone
- 430 VI in LabVIEW.
- 431 To estimate the phase relationship between spikes or membrane currents and the gamma
- 432 oscillation, we first identified successive troughs of the LFP using the WA Multiscale Peak
- 433 Detection VI in LabVIEW. Gamma cycles that deviated more than 20% from the modal
- 434 period were rejected. The membrane current waveform between successive troughs was
- then expressed as a function of instantaneous phase and averaged over all accepted cycles
- 436 in the interval. The minimal (least negative) inward current recorded at -70 mV was
- 437 subtracted to yield the average phasic excitatory current waveform. To estimate the phasic
- 438 inhibitory current waveform, the minimal outward current was subtracted whilst holding cells439 at 0 mV.

440

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#### 445 **Competing interests**

446 The authors declare no competing interests

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#### 652 **Figure Legends**

- Fig. 1. Closed-loop modulation of gamma oscillation 653
- (a) Experimental design. The LFP in CA1 was used to modulate a ramp 654 655 command generated by the PC via a field-programmable gate array (FPGA). 656 The modulated ramp voltage command was then passed to the light-emitting 657 diode (LED) driver, which implemented a threshold-linear voltage-to-current 658 conversion.
- 659 (b) Unmodulated oscillation recorded in CA1 induced by a linear ramp LED driver 660 command. Black trace: LFP with an expanded section showing the 661 characteristic shape of the gamma oscillation (inset). Red trace: LED ramp command. Bottom: LFP Morlet wavelet spectrogram. 662
- 663 (c) Closed-loop oscillation clamp applied between 6 and 8 s, obtained by 664 multiplying the ramp command by  $(1 + k_1 LFP + k_2 dLFP/dt)$ , with dLFP/dtaveraged over 2 ms intervals. For this example,  $k_1 = 0 \text{ mV}^{-1}$ ,  $k_2 = 25 \text{ ms mV}^{-1}$ . 665 The oscillation amplitude was reduced by approximately 60% (insets), with no 666 net change in frequency. 667
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#### Fig. 2. Gamma clamp allows bidirectional modulation of frequency and power

- 671 (a) In-phase modulation led to an increase in gamma power and frequency. Top: 672 200 ms-long segments of the LFP before, during and after closed-loop modulation of the LED driver. Middle: spectrogram. Bottom, left: two cycles of 673 674 the average oscillation before and during the oscillation clamp. The average 675 LED command (red trace, arbitrary scale) is shown superimposed on the 676 clamped oscillation. Right: power spectral density before (blue), during (red) and after (grey) clamp. The polar plot shows the phase relationship between 677 678 the LED command and the LFP. 679
  - (b) Anti-phase modulation led to decreases in both frequency and power.
  - (c) An increase in oscillation frequency, together with a decrease in power, was obtained with ~90° phase-advance of the LED driver command over the LFP.
  - (d) A decrease in frequency, together with an increase in power, was obtained when the LED modulation was delayed relative to the LFP by ~145°. Scale bars apply to all panels.
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#### Fig. 3. Dissociable modulation of oscillation frequency and power

- 687 (a) Dependence of frequency change on the phase relationship between the LED 688 modulation and the LFP (positive values indicate LED phase advance relative 689 to LFP). Changes in frequency are plotted as F<sub>clamped</sub>/F<sub>unclamped</sub>, where the unclamped frequency was averaged from the gamma oscillation for 1 s before 690 and 1 s after the gamma clamp was applied. Data are shown as mean ± SEM 691 692 (n = 19 experiments). A positive phase difference indicates that the 693 modulation was phase-advanced relative to the LFP.
  - (b) Dependence of power change on the phase difference plotted as in (A).
- 695 (c) Change in frequency plotted against change in power for different LED – LFP 696 phase differences (colour code at right).
- (d) Average LFP and phase response (PRC) curve from (Akam et al., 2012) 697 698 (left). The circular cross-correlogram at right yields a prediction of the effect of

a continuous modulation on the oscillation phase, and therefore on its frequency, in rough agreement with the observed relationship in (**a**).

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# 702Fig. 4. Gamma clamp alters the synchrony and phase, rather than rate,703of principal cell firing

- 704 (a) Example closed loop modulation increasing gamma power. Top: red trace 705 showing ramp command. Middle: spectrogram. Bottom: sample traces before 706 (Unclamped) and during (Clamped) closed-loop modulation, showing the LFP 707 and the cell-attached recording with identified spikes highlighted. LFP troughs 708 are indicated by open circles. Six representative LFP traces, aligned by spike time, are shown at right. The polar plot indicates the distribution of spike 709 710 phase for unclamped (black) and clamped (red) periods (averaged from 32 711 trials). The circular histograms sample spikes in 30° bins, and show a 712 decrease in dispersion of spike phase during gamma clamp (LFP trough = 713 0°).
  - (b) Example closed loop modulation increasing gamma frequency, plotted as for (A). The polar plot indicates phase advance of spiking.
- 715 (c) Left: Bidirectional changes in power were associated with corresponding 716 717 changes in the vector length (R) obtained by averaging all spike phases. This 718 is consistent with a decrease in phase dispersion observed with an increase 719 in power, and conversely, an increase in phase scatter with a decrease in 720 power. Changes in power however did not affect the average rate of spiking, 721 when compared with trials when gamma clamp was not applied (middle). 722 Right: increased gamma frequency was associated with a significant phase 723 advance of spiking. \*: p<0.05; \*\*\*: p<0.001. Numbers of experiments are 724 indicated in the bars.
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# 726Fig. 5. Gamma clamp imposes a phasic excitatory current to pyramidal727neurons

- (a) Top: sample LFP (black) and simultaneously recorded holding current in one pyramidal neuron held at -70 mV (blue) before, during and after feedback modulation increasing oscillatory frequency. Bottom: spectrogram.
- (b) Two cycles of the average LFP waveform and membrane current without (Unclamped) and with gamma clamp (Clamped). The average phaseadvanced LED command during feedback modulation is shown superimposed (red). The minimum (least negative) inward current was subtracted (dashed lines) to estimate the phasic excitation. The red arrows indicate the temporal relationship between the peak LED driver command and the maximal excitatory current. Bottom: polar plots indicating the cycleaverage of the excitatory current during unclamped (left) and clamped (right) periods of the trial shown in (a). The vectors indicate the average phases of the currents.
  - (c) Left: difference vector obtained from the vectors in (b), representing the net phasic excitatory current imposed by gamma clamp. Right: phase difference between LED and LFP for the same experiment.
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### Fig. 6. Excitatory current phase determines changes in frequency andpower

- 747(a) Change in gamma frequency and power, plotted against the phase of the net748excitatory current ( $\Delta E$ ) calculated as in Fig. 5. Confidence intervals are SEM749(n = 13).
- 750 **(b)**  $F_{clamped}/F_{unclamped}$  plotted against  $P_{clamped}/P_{unclamped}$  for different excitatory 751 current phases, indicated by the colour code below, aligned with the average 752 LFP waveform. Pyramidal neurons spike around the trough of the LFP (0°).
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## Fig. 7. Gamma clamp afffects phasic inhibitory currents in pyramidal neurons

- (a) Top: sample LFP (black) and simultaneously recorded holding current in one pyramidal neuron held at 0 mV (blue) before, during and after feedback modulation increasing oscillatory frequency. Bottom: spectrogram. Same cell as in Fig. 5.
  - (b) Two cycles of the average LFP waveform and membrane current without (Unclamped) and with gamma clamp (Clamped). The average phaseadvanced LED command during feedback modulation is shown superimposed (red). The minimum (least positive) outward current was subtracted (dashed lines) to estimate the phasic inhibition.
  - (c) Polar plots indicating the cycle-average of the inhibitory current during unclamped (left) and clamped (right) periods of the trial shown in (a). The vectors indicate the average phases of the currents.
- 768 (d) Left: phasic inhibition (vertical scale) plotted against phasic excitation 769 aligning cycle-average (horizontal scale) estimated by membrane 770 conductances by the LFP without (black) and with (red) oscillation clamp with 771 +83° LED-LFP phase advance. Same cell and clamp regime as shown in (a-772 c). Right: representative example of excitatory and inhibitory conductances obtained with 133° LED-LFP phase delay. Same cell but different clamp 773 774 regime. The cycles are arbitrarily anchored to the trough of the LFP (\*). The 775 arrows indicate the direction of the excursion.

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Fig. 2





Fig. 4



### Fig. 5



 $\Delta E$  phase (°)



Fig. 6

